# Toward Direct Determination of Conformations of Protein Building Units from Multidimensional NMR Experiments I. A Theoretical Case Study of For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>

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**ABSTRACT:** NMR chemical shielding anisotropy tensors have been computed, employing several basis sets and the GIAO-RHF and GIAO-MP2 formalisms of electronic structure theory, for all the atoms of the five and nine typical backbone conformers of For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>, respectively. Multidimensional chemical shift plots, as a function of the respective backbone fold, have been generated for both peptide models. On the 2D  ${}^{1}H^{NH}$   ${}^{15}N^{NH}$  and  ${}^{15}N^{NH}$   ${}^{13}C^{\alpha}$  plots the most notable feature is that at all levels of theory studied the backbone conformers cluster in different regions. Computed chemical shifts, as well as their averages, have been compared to relevant experimental values taken from the BioMagnetic Resonance Bank (BMRB). At the highest levels of theory, for all nuclei but the amide protons, deviations between statistically averaged theoretical and experimental shifts are as low as 1-3%. These results indicate that chemical shift information from selected multiple-pulse NMR experiments (e.g., 2D-HSQC and 3D-HNCA) could directly be employed to extract folding information for polypeptides and proteins. © 2000 John Wiley & Sons, Inc. J Comput Chem 21: 882–900, 2000

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Traditionally, determination of the time sional (3D) structure of peptides and small experiments is based on the assignment of proton (<sup>1</sup>H) NMR chemical shifts followed by an analysis of the observed nuclear Overhauser effects (NOEs). The observed NOEs related to spatially close pairs of protons are converted into distances, and these proton-proton distances form the basis of structure analysis.<sup>1-4</sup> Backbone and side-chain dihedral angles can be usually determined using NOE-based constraints even if some of the distances are ambiguous.<sup>4a</sup> For a successful 3D NMR structure determination it is mandatory to assign all resonances. To achieve full assignment for unlabeled and singly labeled (<sup>15</sup>N) proteins, J-correlated spectra, and NOESY-type information are essential. In contrast, if a doubly labeled (13C and <sup>15</sup>N) protein is available, the full assignment can be achieved without NOEs by using specific 3D experiments, exploring only homo- and het-eronuclear coupling constants.<sup>4b-d</sup> Such a strategy could, for example, incorporate the following set of heteronuclear experiments (see ref. 4e for the abbreviations): HNCA, HN(CO)CA, CBCA(CO)NH, HBHA(CO)NH, CC(CO)NH, HCC(CO)NH, and HCCH-TOCSY. At present, however, even if all the assignments are correctly determined from Jcorrelated spectroscopic data, the 3D structure of the molecule cannot be determined without the analysis of the information derived from NOESYtype spectra.

Chemical shielding of a nucleus, located in different proteins or at different sites within the same protein, changes, due either to the individual molecular environment within the macromolecule or to differences in backbone orientations. If the latter factor is dominant, the 3D structure of a protein could be revealed using chemical shift information alone. In the above-mentioned triple-resonance NMR experiments chemical shifts can be resolved since information is spread out, for example, on the  ${}^{1}\text{H}-{}^{15}\text{N}$ , <sup>1</sup>H-<sup>1</sup>H, <sup>1</sup>H-<sup>13</sup>C, and <sup>15</sup>N-<sup>13</sup>C planes. Detailed experimental studies<sup>5-9</sup> have clearly established a few structure-induced  ${}^{13}C^{\alpha}$ ,  ${}^{15}N^{NH}$ , and  ${}^{1}H^{N}$  chemical shift changes in peptides and proteins (vide infra), providing examples of correlation of backbone folds of peptides and proteins with NMR chemical shifts. These results do provide hope that the direct analysis of NMR chemical shifts from relevant multiplepulse experiments (e.g., two-dimensional heteronuclear multiple quantum coherence (2D-HMQC),<sup>10</sup> two-dimensional heteronuclear single quantum coherence (2D-HSQC),<sup>11</sup> and 3D-HNCA<sup>12</sup>) may prove to be a plausible alternative to the distance-based (NOE) strategy for elucidation of the dihedral space of protein structures. The major limitation of this approach, at least at present, is that an unambiguous empirical correlation between backbone conformation and NMR chemical shifts has been put forward only for the  $\alpha$ -helical and  $\beta$ -sheet regions of the Ramachandran surface.<sup>13,14</sup> Due to the lack of appropriate experimental data sets, theoretical (*ab initio*) calculation of conformation-dependent NMR shieldings seems to offer the only hope to establish and probe such relations.<sup>13, 15</sup>

The <sup>33</sup>S chemical shieldings in inorganic and small organic molecules,<sup>16</sup> the <sup>13</sup>C NMR properties<sup>17</sup> of Thr and Tyr amino acids in zwitterions, and the <sup>19</sup>F shifts of various fluorinated tryptophancontaining peptides and proteins<sup>18</sup> have been determined by *ab initio* techniques. Nevertheless, the primary goal of NMR computations on peptides and proteins has been the description of chemical shielding anisotropy (CSA) tensors and associated chemical shifts of the <sup>13</sup>C and <sup>15</sup>N nuclei.<sup>13, 19–24</sup>

Jiao and coworkers<sup>19</sup> were among the first who determined the <sup>13</sup>C CSA tensors of For-Gly-NH<sub>2</sub> as a function of the backbone conformation. The  $C^{\alpha}$ chemical shifts were calculated over the Ramachandran surface using a grid of 30°. It was found that  $^{13}C^{\alpha}$  values in the helical and in the extended regions are shifted by 2.3 ppm downfield and 2.9 ppm upfield, respectively, compared to the random coil value. These shifts should be compared with available experimental results,<sup>19</sup> about +3.2 ppm and -1.2 ppm, respectively. For the alanine residue, Heller and coworkers<sup>23</sup> investigated ways to use  $C^{\alpha}$ CSA tensors for the determination of dihedral angles. Using crosspolarization magic-angle-spinning NMR, the  $C^{\alpha}$  values of tripeptides of known structure were measured. The obtained experimental shifts were correlated with *ab initio* chemical shifts. Using the gauge-including atomic orbital restricted Hartree-Fock (GIAO-RHF) method, the <sup>13</sup>C CSA tensors of the For-L-Ala-NH213c, 20 and For-L-Val-NH2<sup>13c</sup> model compounds were determined, once again as a function of their main-chain fold. It has been concluded that (a) substitution on  $C^{\alpha}$  consistently induces larger shielding shift on  $C^{\beta}$  than on  $C^{\alpha}$ ,<sup>13c</sup> and (b) the  $[\phi, \psi]$  and  $\chi_1$  conformational parameters have a significant influence on chemical shifts. Focusing on Ala and Val residues, Pearson et al.<sup>21</sup> have found that in a protein such as nuclease the  $[\phi, \psi]$  values can be estimated using chemical shifts. Although the results obtained are less dependable than those derived from NOEs and J-coupling constraints, the strategy seemed very promising for the estimation of secondary structures. Analyzing chemical shift nonequivalencies of Ala and Val residues in proteins like calmodulin and nuclease,<sup>13c</sup> the applicability of the direct strategy was further explored.<sup>13c</sup> Prompted by these results, <sup>13</sup>C CSA tensor computations were extended<sup>22</sup> to cover additional amino acid residues Val, Ile, Ser, and Thr. For all residues studied, the  ${}^{13}C^{\alpha}$  shift showed the expected  $\approx 5$  ppm increase for the  $\beta$  conformation over the helical structure. Furthermore, the diagonal CSA tensor elements were all found to be sensitive to changes in the  $\phi$ ,  $\psi$ , and  $\chi_1$  torsion angles. In summary, all calculations indicated that it is possible to deduce certain backbone and sidechain orientations from  ${}^{13}C^{\alpha}$  shielding tensors.

As well known from different <sup>15</sup>N-X-type correlated NMR experiments (e.g., <sup>15</sup>N<sup>NH</sup>-1H<sup>NH</sup> HSQC), <sup>15</sup>N<sup>NH</sup> chemical shifts are rather sensitive probes of protein main-chain fold, and have considerable potential for structure determination. Nevertheless, due to their fairly complex dependence on several torsion angles and on electrostatic field effects,<sup>24</sup> interpretation of changes in <sup>15</sup>N<sup>NH</sup> shifts proved to be somewhat more complicated than that of changes in <sup>13</sup>C shifts. In short, while the most important structural factors determining the <sup>15</sup>N<sup>NH</sup> NMR chemical shift of amino acid residue *i* appear to be the dihedral angles  $\psi^{(i-1)}$  and  $\phi^{(i)}$ , the angles  $\psi^{(i)}$  and  $\phi^{(i-1)}$ and the side-chain orientation also have a nonnegligible effect. Consequently, at present, <sup>15</sup>N<sup>NH</sup> shifts are less readily useful than <sup>13</sup>C shifts.

H-bonds have a crucial role in the formation and stabilization of peptides and proteins. For example, Asakawa and coworkers<sup>25</sup> have found that the  $C^{\alpha}$  shielding values are affected not only by the values of the  $[\phi, \psi]$  torsional parameters but also by H-bond patterns (standard, bifurcated, etc.) and by H-bond strength. This observation proved to be crucial during determination of the 3D structure of ribonuclease and that of a basic pancreatic trypsin inhibitor.<sup>25</sup> GIAO-RHF calculations also established the influence of H-bonding on the carbonyl carbon of N-methylacetamide interacting with formamide.26 Further computations,26 at the 6-31G\*\* GIAO-RHF level, have been performed on the helical and  $\beta$ -sheet structures of For-(Ala)<sub>5</sub>-NH<sub>2</sub> in order to investigate the effect of intramolecular H-bonds on shieldings of different nuclei, for example, <sup>13</sup>C'. The experimentally determined difference between the helical and the  $\beta$ -sheet values of  ${}^{13}C'$ , 4.6 ppm, is close to the calculated value, 4.9 ppm. The most significant perturbation caused by the H-bond was on carbonyl carbons.<sup>27</sup>

When performing any type of ab initio computation on any property of a chemical system, issues related to the applied theoretical level (n-particle space) and basis set (one-particle space) are to be considered. Laws and coworkers<sup>28</sup> have determined, at the GIAO-RHF level, the effect of basis set extension, employing standard basis sets STO-3G, 3-21G, 4-31G, and 6-311+G(2d), on  ${}^{13}C^{\alpha}$  and  ${}^{13}C^{\beta}$  shifts of two amino acid residues, Ala and Val, in proteins. They found that the chemical shifts of these residues, determined using smaller basis sets, correlate rather well with results obtained with larger basis sets, even when the underlying reference geometries are slightly different. It is encouraging that chemical shifts computed at different levels can be scaled and fine tuned. In a further study of relevance, employing empirical chemical shift surfaces, an attempt was made to predict <sup>13</sup>C shift values of valine residues in three well-known proteins (calmoduline, nuclease, and ubiquitine)<sup>29</sup> utilizing their X-ray structures. Most HF results and experimental values showed agreement<sup>29</sup> poorer than expected, but the agreement improved slightly by using density functional theory (DFT)-type calculations. Pearson and coworkers<sup>29</sup> have thus concluded that for accurate computation of chemical shifts geometry optimization and the inclusion of electron correlation in the theoretical treatment appear to be important.

In this report a concerted attempt is made to correlate calculated isotropic NMR shielding and chemical shift values with all characteristic backbone conformations<sup>30</sup> of peptide models, concentrating not only on <sup>13</sup>C but also on <sup>15</sup>N and <sup>1</sup>H. The effect of electron correlation is explicitly incorporated in CSA tensor calculations at the GIAO-MP2 level. To keep the computational problem tractable, systematic calculations are carried out only for conformations of the following amino acid diamides of glycine and alanine: For-Gly-NH2 and For-L-Ala-NH<sub>2</sub>. It is not obvious what is the best use of the large number of quantum chemical data obtained. We felt that apart from investigating traditional issues, such as basis set dependence and the effect of electron correlation on the calculated results, a thorough statistical analysis offers the best insight. Therefore, relevant correlations are investigated in considerable detail, confirming some well-known structure-chemical shift relations and establishing, at the same time, new and potentially useful ones. Although detailed comparison of the computed results with experiment is not yet possible, an attempt is made to compare direct and statistically averaged chemical shifts from the two sources.

#### **Computational Details**

Computation of NMR shielding tensors utilized the Gaussian94,<sup>31</sup> Gaussian98,<sup>32</sup> and AcesII<sup>33</sup> program systems, and were performed at the GIAO-RHF (gauge including atomic orbitals restricted Hartree–Fock)<sup>34</sup> and GIAO-MP2 (GIAO secondorder Møller–Plesset perturbation theory)<sup>35</sup> levels of theory employing basis sets of  $6-31+G^{*,36a}$  6- $311++G^{**,36b}$  DZ(d),<sup>37</sup> and TZ2P<sup>38</sup> quality.

The For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub> models are depicted on Figures 1 and 2, respectively. These figures also contain information concerning average chemical shifts of all nuclei forming the amino acid core, as taken from the BioMagnetic Resonance Bank (BMRB).<sup>39</sup>

The reference geometries employed for the chemical shielding anisotropy (CSA) computations have been determined at the 6-31++G\*\* RHF and 6-311++G\*\* DFT(B3LYP) levels. The latter level was employed extensively for the geometry optimizations, as it proved to be especially reliable during structural studies of the neutral amino acids glycine<sup>40a, 40b</sup> and  $\alpha$ -alanine.<sup>40c</sup> To obtain a complete chemical shielding set associated with the five and nine unique backbone orientations of For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>, respectively, additional 6-31++G\*\* RHF constrained geometry optimizations have been performed keeping the characteristic  $\phi$  and  $\psi$  torsional angles constant. These fixed  $\phi$  and  $\psi$  values were obtained from averaging the appropriate 3-21G RHF backbone conformational parameters of For-L-Ala-L-Ala-NH2.44 Description of Levels A-F, resulting from different combinations of full/constrained geometry optimizations and CSA calculations, is given in Table I.



**FIGURE 1.** The For-Gly-NH<sub>2</sub> model and average chemical shifts (in ppm) taken from BMRB.



**FIGURE 2.** The For-L-Ala-NH<sub>2</sub> model and average chemical shifts (in ppm) taken from BMRB.

Selected geometric parameters (including optimized/fixed  $\phi$  and  $\psi$  values), energies, and isotropic chemical shielding values ( $\sigma$ -scale) are reported for For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub> in Tables II–VI. When relative chemical shifts ( $\delta$ -scale) are used, the appropriate isotropic chemical shielding values of <sup>1</sup>H, <sup>13</sup>C, and <sup>15</sup>N are referenced to <sup>1</sup>H and <sup>13</sup>C of tetramethyl-silane (TMS) and to <sup>15</sup>N of NH<sub>3</sub>. The reference geometry chosen for NH<sub>3</sub> corresponds to the aug-cc-pVTZ CCSD(T) optimized geometry,<sup>42</sup> while the geometry of TMS has been optimized at the 6-311++G<sup>\*\*</sup> B3LYP level (cf. Tables V and VI).

Average (or random) chemical shift data and associated standard deviations for the amino acid residues Gly and Ala were taken from data deposited in BMRB<sup>39</sup> (cf. Figs. 1 and 2 and Tables V and VI). Experimental chemical shifts associated with each of the 9(5) typical backbone conformers of For-L-Ala-NH<sub>2</sub>(For-Gly-NH<sub>2</sub>) were determined by a comprehensive analysis of data deposited in BMRB and the Protein Data Bank (PDB).<sup>43</sup>

#### **Structures and Energetics**

The usual way to characterize the backbone conformation of peptides (or proteins) is through the  $[\phi, \psi]$  torsional angle pairs associated with the amino acid residues.<sup>30</sup> The torsional potentials along both  $\phi$  and  $\psi$  should have three minima; therefore, in this simplest picture we expect nine legitimate conformers for each "peptide unit" (...NH-CHR-CO...).<sup>30, 41, 44</sup> These nine conformers are named, according to our established convention,<sup>30c</sup> as follows:  $\alpha_L, \alpha_D, \beta_L, \gamma_L, \gamma_D, \delta_L, \delta_D, \varepsilon_L$ , and  $\varepsilon_D$  (see Figs. 3 and 4 for the Ramachandran surfaces of For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>, respectively). The lack of configurational chirality of

#### TABLE I.

For-Gly-NH <sub>2</sub> a	nd For-	L-Ala-NH <sub>2</sub> Model System	ıs.		
Model System	Level	NMR Computation	Geometry Optimization	No. of Structures	Comment
For-Gly-NH <sub>2</sub>	А	GIAO-RHF/DZ(d)	B3LYP/6-311++G**	5	full opt. plus $[\phi, \psi]$ constr
	В	GIAO-MP2/DZ(d)	B3LYP/6-311++G**	5	full opt. plus $[\phi, \psi]$ constr
	С	GIAO-RHF/TZ2P	B3LYP/6-311++G**	5	full opt. plus $[\phi,\psi]$ constr
For-L-Ala-NH <sub>2</sub>	А	GIAO-RHF/6-31+G*	RHF/6-31++G**	9	$[\phi,\psi]$ constr.
	В	GIAO-RHF/6-31+G*	B3LYP/6-311++G**	6	full opt. only
	С	GIAO-RHF/6-31+G*	RHF/6-31++G**	6	full opt. only
	D	GIAO-RHF/6-311++G**	RHF/6-31++G**	9	$[\phi, \psi]$ constr.
	Е	GIAO-RHF/6-311++G**	B3LYP/6-311++G**	6	full opt. only
	F	GIAO-RHF/TZ2P	B3LYP/6-311++G**	6	full opt. only

B3LYP/6-311++G\*\*

Designation of Different Combinations of Geometry Optimization and Chemical Shielding Calculations for the For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub> Model Systems.

glycine makes it unique among the natural alpha amino acids resulting in a double degeneracy on the Ramachandran PES. In our notation this means that for For-Gly-NH<sub>2</sub>  $\alpha_L = \alpha_D = \alpha_{LD}$ ,  $\beta_L = \beta_{LD}$ ,  $\gamma_L = \gamma_D = \gamma_{LD}$ ,  $\varepsilon_L = \varepsilon_D = \varepsilon_{LD}$ , and  $\delta_L = \delta_D = \delta_{LD}$ . Of course, if the glycine residue is embedded in a chiral environment (i.e., in a protein chain), the de-

GIAO-RHF/TZ2P

G

generacy of the conformer pairs of our For-Gly- $NH_2$  model disappears.

full opt. and  $[\phi, \psi]$  constr.

10

Although classification of hundreds of X-ray determined protein structures revealed the existence of all nine conformer types, systematic *ab initio* calculations<sup>30, 45</sup> carried out on simple amino acid diamides (For-X-NH<sub>2</sub>, with X = Gly, Ala, Val, Ser,

TABLE II.

Conformational Parameters (Torsion Angles  $\phi$  and  $\psi$  in Degrees), Total Energies ( $E/E_h$ ), Relative Energies [ $\Delta E/(kcal mol^{-1})$ ], and Relative Populations of the Characteristic Conformers of For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>.<sup>a</sup>

Model	Backbone Conformation	$\phi$	$\psi$	E/E <sub>h</sub>	$\Delta E/(\text{kcal mol}^{-1})$	$p(i)/\sum p(i)$
For-Gly-NH <sub>2</sub>	$\alpha_{\rm ID}$	-68.6 <sup>b</sup>	-17.5 <sup>b</sup>	-378.027433	3.21	0.002
• -	$\beta_{LD}^{}$	180.00	180.00	-378.032393	0.10	0.444
	γld	-81.36	65.78	-378.032551	0.00	0.524
	$\delta_{LD}$	113.18	-18.26	-378.029586	1.86	0.023
	$\varepsilon_{LD}$	-74.7 <sup>c</sup>	167.8 <sup>c</sup>	-378.028406	2.60	0.007
For-L-Ala-NH <sub>2</sub>	$\alpha_{D}$	73.4 <sup>b</sup>	10.5 <sup>b</sup>	-417.352003	5.00	0.000
	$\alpha_{L}$	-68.6 <sup>b</sup>	-17.5 <sup>b</sup>	-417.354165	3.65	0.002
	$\beta_{L}$	-156.59	163.69	-417.358625	0.85	0.185
	γD	72.48	-55.06	-417.356591	2.12	0.022
	γL	-82.60	73.44	-417.359976	0.00	0.770
	$\delta_{D}$	-168.06	-42.03	-417.350367	6.03	0.000
	$\delta_{L}$	-114.02	15.76	-417.356082	2.44	0.013
	₽D	60 <sup>d</sup>	-120 <sup>d</sup>	-417.351456	5.35	0.000
	εL	-60 <sup>d</sup>	120 <sup>d</sup>	-417.355659	2.71	0.008

<sup>a</sup> All results reported, unless otherwise noted, correspond to the 6-311++G<sup>\*\*</sup> B3LYP level. Relative populations,  $p(i) / \sum p(i)$ , for computed *ab initio* relative energies are determined as  $exp(-\Delta E/RT) / \sum exp(-\Delta E/RT)$ , where RT = NkT = 0.595371 kcal mol<sup>-1</sup>, T = 300 K,  $k = 1.38 \times 10^{-23}$  J/K, and Avogadro's number (N) is 6.02  $\times 10^{23}$  mol<sup>-1</sup>.

<sup>b</sup> Constrained backbone parameters of helical conformers found for For-L-Ala-L-Ala- $H_2$ .<sup>41</sup>

<sup>c</sup> Constrained backbone parameters of polyproline II conformers found for For-L-Ala-L-Ala-NH<sub>2</sub>.<sup>41</sup>

<sup>d</sup> For symmetry reasons these idealized backbone values were used, which are characteristic for polyproline II.

#### TABLE III.

## Backbone Conformational Parameters and Isotropic Chemical Shielding Values ( $\sigma$ -Scale), Relative to the $\gamma_{LD}$ Conformation, of For-Gly-NH<sub>2</sub>.<sup>a</sup>

Backbone Conformation	Level <sup>b</sup>	ω	$\phi$	$\psi$	<i>ω</i> 1	<sup>15</sup> N <sup>H</sup>	<sup>13</sup> C <sup>α</sup>	<sup>13</sup> C′	<sup>1</sup> H <sup>N</sup>	$^{1}$ H $^{lpha}$ (Average)
$\alpha_{\rm LD}$	A	-175.4	-68.6	-17.5	173.1	11.40	3.68	4.86	1.09	0.32
	B	-175.4	-68.6	-17.5	173.1	11.32	4.85	3.13	1.13	0.35
	C	-175.4	-68.6	-17.5	173.1	4.35	1.00	–0.15	0.35	-0.01
$\beta_{LD}$	A	180.0	180.0	180.0	180.0	7.97	4.73	0.37	-0.67	0.04
	B	180.0	180.0	180.0	180.0	8.67	6.42	0.87	-0.18	-0.01
	C	180.0	180.0	180.0	180.0	10.03	5.02	0.00	-1.02	-0.08
γld	A	-177.6	-81.4	65.8	-178.2	0.00	0.00	0.00	0.00	0.00
	B	-177.6	-81.4	65.8	-178.2	0.00	0.00	0.00	0.00	0.00
	C	-177.6	-81.4	65.8	-178.2	0.00	0.00	0.00	0.00	0.00
$\delta_{LD}$	A	173.5	113.2	-18.3	-172.7	1.67	3.54	-0.12	0.63	-0.27
	B	173.5	113.2	-18.3	-172.7	2.13	5.46	0.45	0.62	-0.19
	C	173.5	113.2	-18.3	-172.7	1.81	3.66	-0.25	0.64	-0.40
€LD	A	167.1	-74.7	167.8	-175.3	13.28	6.08	4.07	1.00	0.53
	B	167.1	-74.7	167.8	-175.3	13.94	8.01	1.96	1.07	0.51
	C	167.1	-74.7	167.8	-175.3	7.45	3.91	-1.40	0.40	0.09

<sup>a</sup> All relative isotropic shielding values (in ppm) are referenced to the absolute isotropic shieldings of the  $\gamma_{LD}$  conformer, which are as follows: Level A:  ${}^{15}N^{NH} = 164.7$ ,  ${}^{13}C^{\alpha} = 155.5$ ,  ${}^{13}C' = 26.9$ ,  ${}^{1}H^{NH} = 28.5$ ,  ${}^{1}H^{\alpha} = 29.2$ ; Level B:  ${}^{15}N^{NH} = 168.3$ ,  ${}^{13}C^{\alpha} = 157.3$ ,  ${}^{13}C' = 50.1$ ,  ${}^{1}H^{NH} = 27.7$ ,  ${}^{1}H^{\alpha} = 28.6$ ; Level C:  ${}^{15}N^{NH} = 151.8$ ,  ${}^{13}C^{\alpha} = 149.2$ ,  ${}^{13}C' = 14.5$ ,  ${}^{1}H^{NH} = 27.0$ ,  ${}^{1}H^{\alpha} = 28.5$ . <sup>b</sup> See Table I for the designation of the *ab initio* levels employed.

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Backbone Conformation	Level <sup>b</sup>	ω <sub>0</sub>	$\phi$	$\psi$	ω <sub>1</sub>	χ1	<sup>15</sup> N <sup>NH</sup>	<sup>13</sup> C <sup>α</sup>	<sup>13</sup> C <sup>β</sup>	<sup>13</sup> C′	<sup>1</sup> H <sup>NH</sup>	<sup>1</sup> Η <sup>α</sup>
α <sub>D</sub>	A	171.9	61.8	31.9	-173.9	64.2	8.32	-5.27	-0.28	0.06	-0.12	0.89
	B	171.8	73.4	10.5	-173.6	61.9	8.51	-6.49	0.03	-0.62	-0.03	0.88
	C	169.5	69.9	24.6	-174.7	64.8	8.37	-5.80	0.42	-0.04	-0.02	0.91
	D	171.9	61.8	31.9	-173.9	64.2	9.31	-5.74	-0.36	-0.19	-0.17	0.94
	E	171.8	73.4	10.5	-173.6	61.9	9.56	-6.91	-0.15	-1.02	-0.10	0.90
	F	171.8	73.4	10.5	-173.6	61.9	10.02	-6.87	0.08	-1.34	-0.03	0.97
αL	A	-175.3	-68.6	-17.5	172.9	61.7	3.97	-4.11	-1.62	-1.63	0.17	0.17
	D	-175.3	-68.6	-17.5	172.9	61.7	4.14	-4.14	-1.98	-1.83	0.08	0.17
$3_{10}$ -helix like <sup>c</sup>	F	-176.3	-68.6	-17.5	172.7	-59.9	4.68	-3.74	-2.03	-1.79	0.21	0.11
$\alpha$ -helix-like	F	-175.7	-54.0	-45.0	173.1	61.1	6.25	-5.50	-2.27	-1.83	0.32	0.25
βL	A	-179.6	-167.6	169.9	-179.1	62.4	10.26	-1.53	-3.07	-0.81	-1.45	0.20
	B	177.2	-156.6	163.7	178.7	61.8	9.81	-0.51	-4.42	-1.04	-1.49	0.10
	C	177.9	-155.6	160.4	178.7	61.4	8.81	-1.30	-3.98	-1.00	-1.29	0.11
	D	-179.6	-167.6	169.9	-179.1	62.4	10.91	-1.65	-3.41	-0.87	-1.63	0.17
	E	177.2	-156.6	163.7	178.7	61.8	10.60	-0.54	-4.74	-1.45	-1.65	0.05
	F	177.2	-156.6	163.7	178.7	61.8	10.75	-0.41	-4.55	-1.81	-1.69	0.07

#### TABLE IV. \_\_\_\_

(Co	ntinu	ıed)
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Backbone Conformation	Level <sup>b</sup>	ω	$\phi$	$\psi$	$\omega_1$	χ1	<sup>15</sup> N <sup>NH</sup>	<sup>13</sup> C <sup>α</sup>	<sup>13</sup> C <sup>β</sup>	<sup>13</sup> C′	<sup>1</sup> H <sup>NH</sup>	<sup>1</sup> Ηα
γD	A	177.6	74.3	-59.5	-179.8	63.4	7.30	-10.59	-0.21	-0.91	-0.30	0.73
	B	175.8	72.5	-55.1	-178.4	62.7	5.94	-10.35	-0.87	-0.83	-0.26	0.73
	C	176.4	75.2	-55.0	-177.8	63.7	6.44	-10.67	-0.14	-1.18	-0.23	0.74
	D	177.6	74.3	-59.5	-179.8	63.4	7.98	-11.25	-0.80	-1.14	-0.37	0.80
	E	175.8	72.5	-55.1	-178.4	62.7	6.68	-11.02	-1.47	-1.30	-0.32	0.78
	F	175.8	72.5	-55.1	-178.4	62.7	6.31	-11.14	-1.35	-1.20	-0.30	0.83
γL	A	-178.0	-84.5	68.7	-176.9	62.1	0.00	0.00	0.00	0.00	0.00	0.00
	B	-177.5	-82.6	73.4	-173.8	63.8	0.00	0.00	0.00	0.00	0.00	0.00
	C	-177.8	-86.2	76.5	-173.6	62.1	0.00	0.00	0.00	0.00	0.00	0.00
	D	-178.0	-84.5	68.7	-176.9	62.1	0.00	0.00	0.00	0.00	0.00	0.00
	E	-177.5	-82.6	73.4	-173.8	63.8	0.00	0.00	0.00	0.00	0.00	0.00
	F	-177.5	-82.6	73.4	-173.8	63.8	0.00	0.00	0.00	0.00	0.00	0.00
δ <sub>D</sub>	A	173.5	-179.6	-43.7	-170.1	55.4	9.65	-8.08	-1.85	-1.32	0.19	0.12
	B	172.3	-168.1	-42.0	-170.9	58.7	10.39	-6.66	-2.57	-1.05	0.14	-0.01
	C	171.4	-165.9	-41.9	-170.1	58.2	9.70	-7.38	-2.20	-1.33	0.09	0.04
	D	173.5	-179.6	-43.7	-170.1	55.4	10.75	-8.72	-2.55	-1.56	0.18	0.14
	E	172.3	-168.1	-42.0	-170.9	58.7	11.39	-7.32	-3.17	-1.52	0.12	-0.02
	F	172.3	-168.1	-42.0	-170.9	58.7	11.29	-7.30	-3.15	-1.70	0.08	-0.04
δ	A	-170.5	-126.2	26.5	171.0	61.0	2.22	-0.30	-2.52	-0.66	0.36	-0.54
	B	-171.5	-114.0	15.8	171.8	59.4	2.56	-0.31	-2.57	-0.54	0.46	-0.50
	C	-171.0	-111.8	13.5	171.2	60.6	2.27	-0.94	-2.45	-1.26	0.43	-0.41
	D	-170.5	-126.2	26.5	171.0	61.0	2.46	-0.41	-2.56	-0.90	0.32	-0.62
	E	-171.5	-114.0	15.8	171.8	59.4	2.73	-0.43	-2.69	-0.69	0.40	-0.60
	F	-171.5	-114.0	15.8	171.8	59.4	2.77	-0.40	-2.70	-0.97	0.49	-0.58
<sup>2</sup> D	A	-163.9	64.7	-178.6	176.2	73.1	12.65	-5.48	-1.44	0.22	0.20	0.93
	D	-163.9	64.7	-178.6	176.2	73.1	13.91	-6.18	-1.73	0.04	0.13	1.00
	F	-173.2	60.0	-120.0	172.3	60.3	8.11	-5.56	-0.70	-3.17	-0.56	1.02
εL	A	166.9	-74.7	167.8	-174.0	62.1	6.73	-1.36	-2.31	-2.50	0.04	0.02
	D	166.9	-74.7	167.8	-174.0	62.1	7.35	-1.36	-2.84	-2.81	0.01	0.00
	F	171.2	-60.0	120.0	-173.1	64.0	3.48	-2.34	-0.78	-5.54	0.02	0.48

<sup>a</sup> All relative isotropic shielding values are referenced to the absolute isotropic shielding of the  $\gamma_{L}(g+)$  conformer calculated as: Level A:  ${}^{15}N^{NH} = 158.7$ ,  ${}^{13}C^{\alpha} = 159.4$ ,  ${}^{13}C^{\beta} = 185.9$ ,  ${}^{13}C' = 33.5$ ,  ${}^{1}H^{NH} = 28.9$ ,  ${}^{1}H^{\alpha} = 28.8$ ; Level B:  ${}^{15}N^{NH} = 151.6$ ,  ${}^{13}C^{\alpha} = 156.4$ ,  ${}^{13}C^{\beta} = 184.9$ ,  ${}^{13}C' = 28.3$ ,  ${}^{1}H^{NH} = 28.5$ ; Level C:  ${}^{15}N^{NH} = 159.2$ ,  ${}^{13}C^{\alpha} = 159.5$ ,  ${}^{13}C^{\beta} = 185.9$ ,  ${}^{13}C' = 33.4$ ,  ${}^{1}H^{NH} = 28.9$ ,  ${}^{1}H^{\alpha} = 28.8$ ; Level D:  ${}^{15}N^{NH} = 141.5$ ,  ${}^{13}C^{\alpha} = 151.9$ ,  ${}^{13}C^{\beta} = 179.4$ ,  ${}^{13}C' = 17.9$ ,  ${}^{1}H^{NH} = 28.3$ ,  ${}^{1}H^{\alpha} = 28.5$ ; Level E:  ${}^{15}N^{NH} = 134.1$ ,  ${}^{13}C^{\alpha} = 148.8$ ,  ${}^{13}C^{\beta} = 178.3$ ,  ${}^{13}C' = 12.8$ ,  ${}^{1}H^{NH} = 27.7$ ,  ${}^{1}H^{\alpha} = 28.2$ ; Level F:  ${}^{15}N^{NH} = 136.1$ ,  ${}^{13}C^{\alpha} = 148.6$ ,  ${}^{13}C^{\beta} = 178.1$ ,  ${}^{13}C' = 12.5$ ,  ${}^{1}H^{NH} = 27.4$ ,  ${}^{1}H^{\alpha} = 28.0$ .

<sup>b</sup> See Table I for the designation of the *ab initio* levels employed.

and Phe) resulted in fewer minima than the maximum number allowed. For example, at the 3-21G(6-311++G<sup>\*\*</sup>) RHF level two(three) conformers could not be located for For-L-Ala-NH<sub>2</sub>. The results of these systematic calculations<sup>30, 44, 45</sup> show that the  $\alpha_{\rm L}$ ,  $\varepsilon_{\rm L}$ , and  $\varepsilon_{\rm D}$  conformers are found to be absent most often. As far as relative energies are concerned (see Table II), the computational results reveal that the  $\alpha_{\rm L}$  and  $\varepsilon_{\rm L}$  conformations have high relative energies compared to the usually most stable  $\gamma_L$  conformation. Nevertheless, existence of the "missing" minima has been verified in other amino acid diamides (e.g., For-L-Ser-NH<sub>2</sub>)<sup>44a</sup> as well as in larger peptide models (e.g., For-L-Ala-L-Ala-NH<sub>2</sub>).<sup>41</sup>

It is not the principal purpose of this article to discuss structural and energetic features of the two investigated model compounds at considerable length. Nevertheless, it is worth discussing the re-

		Avera	ige Shifts (pp	om)			I	Error (%)	d	
Level <sup>c</sup>	<sup>15</sup> N <sup>NH</sup> avr	$^{13}C^{lpha}_{avr}$	$^{13}C_{avr}^{\prime}$	$^{1}H_{avr}^{NH}$	${}^{1}H^{\alpha}_{avr}$	Ν	$C^{lpha}$	C′	NH	Hα
Arithmetic Avera	ge									
А	90.97	40.49	170.84	3.88	3.50	17	10	2	53	10
В	105.34	44.06	154.86	3.83	3.35	3	3	11	54	14
С	112.21	42.78	180.50	5.20	3.80	3	5	4	28	3
Boltzmann Avera	age <sup>f</sup>									
А	88.35	39.86	169.91	4.14	3.45	19	12	2	50	11
В	108.53	45.96	155.73	4.42	3.48	1	2	10	47	11
С	112.38	43.15	180.16	5.71	3.76	3	5	4	32	3
Expt.	109.1	45.24	173.84	8.33	3.90					

#### TABLE V. \_\_\_\_\_ Chemical Shifts ( $\delta$ -Scale) Averaged over All Backbone Conformers of For-Gly-NH<sub>2</sub>.<sup>a,b</sup>

<sup>a</sup> Absolute chemical shifts (in ppm) used as references are given as follows: Level A: <sup>1</sup>H (TMS) = 32.800, <sup>13</sup>C (TMS) = 199.590, <sup>15</sup>N (NH<sub>3</sub>) = 262.490; Level B: <sup>1</sup>H (TMS) = 32.070, <sup>13</sup>C (TMS) = 206.265, <sup>15</sup>N(NH<sub>3</sub>) = 280.869; Level C: <sup>1</sup>H (TMS) = 32.231, <sup>13</sup>C (TMS) = 194.688, <sup>15</sup>N (NH<sub>3</sub>) = 268.698.

<sup>b</sup> The number of backbone conformers is 5.

<sup>c</sup> See Table I for the designation of the *ab initio* levels used. Expt. = average experimental shift taken form BMRB.<sup>39</sup>

<sup>d</sup> Error =  $abs(CSA^{exp} - CSA^{calc})/CSA^{exp}*100$ .

<sup>e</sup> Standard deviation of the average shift.

<sup>f</sup> Boltzmann averaging over all conformers was achieved with the use of the corresponding Level A-F energies of the GIAO calculations.

sults obtained and compare them with previously determined values.<sup>46</sup>

The  $\phi$  and  $\psi$  torsional angles do not change much among the model compounds For-Gly-NH<sub>2</sub>, For-L-Ala-NH<sub>2</sub>, and For-L-Ala-L-Ala-NH<sub>2</sub>. For example, the constrained  $\phi$  and  $\psi$  torsion angles taken<sup>41</sup> from For-L-Ala-L-Ala-NH<sub>2</sub> never deviate more than 15° from the optimized ones (cf. Table II). The structural characteristics, when appropriate, are similar to those observed for the parent amino acids, glycine<sup>40a</sup> and  $\alpha$ -alanine.<sup>40c</sup>

The relative energies of the conformers presented in Table II reveal that the energy order of the L-type conformers (cf. Figs. 3 and 4) is identical for both models, regardless of whether electron correlation is incorporated into the computation or not (relevant energy data, not detailed here, can be obtained from the authors upon request). The observed energy order of the L-type conformers is as follows:  $E(\gamma_{\rm L}) < E(\beta_{\rm L}) < E(\delta_{\rm L}) < E(\varepsilon_{\rm L}) < E(\alpha_{\rm L})$ . Typically, the D-type backbone conformers of L-amino acid diamides have a lower stability. The most stable D-type conformer,  $\gamma_D$  of For-L-Ala-NH<sub>2</sub>, has a low relative energy due to a strong intramolecular H-bond, which makes its relative energy similar to that of  $\delta_{\rm L}$ . The  $\delta_{\rm D}$  conformer, though a minimum on the PES of For-L-Ala-NH<sub>2</sub>, has a very high relative energy. The  $\alpha_D$  conformer of For-L-Ala-NH<sub>2</sub>,

a minimum at the 6-311++G<sup>\*\*</sup> RHF level, is not a stationary point at the 6-311++G<sup>\*\*</sup> DFT(B3LYP) level. For For-Gly-NH<sub>2</sub>, the RHF energy difference between the  $\gamma_{\text{LD}}$  and  $\beta_{\text{LD}}$  conformers is typically 1 kcal mol<sup>-1</sup>, favoring  $\gamma_{\text{LD}}$ . However, electron correlation stabilizes  $\beta_{\text{LD}}$  over  $\gamma_{\text{LD}}$ . As a result, while the RHF energy order,  $E(\gamma_{\text{L}}) < E(\beta_{\text{L}})$ , is preserved, using the DFT(B3LYP) or MP2 methods, the difference between the total energies of the two conformers becomes almost negligible.

#### **Accuracy of Computed Chemical Shifts**

Average NMR chemical shifts (and their standard deviations) of all nuclei of glycine and alanine residues are available from BMRB.<sup>39</sup> These data are reported in Tables V and VI. Therefore, one can attempt to answer how well the computed chemical shifts compare with their experimental counterparts. The experimental shifts of the different nuclei ( $^{15}N^{NH}$ ,  $^{13}C^{\alpha}$ ,  $^{13}C^{\beta}$ ,  $^{13}C'$ ,  $^{1}H^{NH}$ ) retrieved from the database refer to a hypothetical average backbone conformation of the relevant amino acid residue. For purposes of comparison either a simple arithmetical average or a Boltzmann-type (energyweighted) average of the shifts of the nuclei of individual conformers can be determined. These

I ABLE Chemic	VI al Shifts (ô-Sc	ale) Averaged ove	r All Backbone C	Conformers of F	-or-L-Ala-NH <sub>2</sub> . <sup>a,b</sup>								
			Av	rerage Chemica	l Shifts (ppm)					Error	p(%)		
Level <sup>c</sup>	No. Conf.	15NNH avr	$^{13}C^{lpha}_{avr}$	$^{13}C^{\beta}_{avr}$	<sup>13</sup> C <sup>′</sup> <sub>avr</sub>	<sup>1</sup> H <sup>NH</sup> avr	$^{1}H_{avr}^{lpha}$	z	ŭ	Cβ	ò	H	Ηα
Arithmeti	cal Average												
۷	, 6	101.46 (4.06) <sup>e</sup>	45.91 (3.65)	16.81 (1.10)	168.58 (0.89)	3.97 (0.54)	3.69 (0.48)	17	13	11	Ŋ	51	13
В	9	109.09 (4.19)	48.89 (4.37)	18.05 (1.76)	173.64 (0.39)	4.73 (0.68)	4.09 (0.52)	1	ø	ß	2	42	4
U	9	101.80 (3.93)	46.06 (4.27)	16.74 (1.75)	168.64 (0.62)	4.09 (0.59)	3.75 (0.50)	17	13	11	Ŋ	50	42
Δ	6	113.86 (4.47)	48.04 (3.92)	17.97 (1.18)	178.69 (0.94)	4.25 (0.59)	3.62 (0.53)	7	6	ß	-	48	15
ш	9	121.82 (4.60)	51.09 (4.66)	19.26 (1.85)	183.78 (0.58)	5.02 (0.72)	4.04 (0.56)	-	4	2	4	39	ß
ш	9	125.76 (4.67)	50.49 (4.72)	18.92 (2.16)	183.34 (0.65)	5.12 (0.75)	4.04 (0.59)	2	ß	0	ო	37	ß
Boltzmar	าท Average <sup>f</sup>												
A	6	104.77	42.46	16.38	168.06	4.33	3.91	15	20	13	ß	47	ω
В	9	110.45	45.16	18.51	173.48	5.25	4.24	10	15	N	N	36	0
U	9	104.01	42.34	17.03	168.28	4.45	3.93	15	20	10	ß	46	~
Δ	0	117.11	44.41	17.50	178.06	4.67	3.85	Ŋ	16	7	0	43	ი
ш	9	122.97	47.08	19.77	183.56	5.63	4.20	0	11	ß	ო	31	-
ш	9	126.50	46.43	19.26	183.20	5.83	4.21	ო	12	2	ო	29	-
Expt.		122.82 (5.88)	53.05 (2.17)	18.91 (1.78)	177.40 (8.32)	8.17 (0.62)	4.24 (0.47)						
<sup>a</sup> Absc 266.95 194.68	lute chemical st 7; Levels D and 8, <sup>15</sup> N (NH <sub>3</sub> ) =	nifts (in ppm) used as E (6-311++G** GIAC 268.698.	s references are give -RHF): <sup>1</sup> H (TMS) =	en as follows: Lev : 32.415, <sup>13</sup> C (TMS	els A, B, and C (6-3 5) = 195.559, <sup>15</sup> N (h	:1+G* GIAO-RH 1H <sub>3</sub> ) = 262.784;	F): <sup>1</sup> H (TMS) = 3 Level F (TZ2P Gl	2.000, <sup>1</sup> AO-RHF	<sup>3</sup> C (TM: 1: <sup>1</sup> H (TI	S) = 20 MS) = 3	1.239, 32.231,	<sup>15</sup> N (NH <sup>13</sup> C (TM	3) =   
<sup>b</sup> The I <sup>c</sup> See <sup>-</sup> <sup>d</sup> Error	number of backt fable I for the de = $abs(\delta^{exp} - \delta^{t})$	oone conformers is 9 signation of the <i>ab in</i> calc)/δexp∗100.	when constrained $\phi$ 't <i>io</i> levels used. Expl	, $\psi$ and 6 when o t. = average expe	ptimized $\phi,\psi$ value rimental shift taken i	s are used. rom BMRB. <sup>39</sup>							

## PERCZEL AND CSÁSZÁR

<sup>e</sup> Standard deviation of the average shift is given in parentheses. <sup>f</sup> Boltzmann averaging over all conformers was achieved with the use of the corresponding Level A–F energies of the GIAO calculations.



**FIGURE 3.** The 3-21G RHF, 6-31+G\* RHF, 6-311++G\*\* B3LYP and the idealized locations of the different backbone conformers of For-Gly-NH<sub>2</sub> on a Ramachandran surface. Note that the lack of configuration chirality results in  $\alpha_L = \alpha_D = \alpha_{LD}$ ,  $\beta_L = \beta_{LD}$ ,  $\gamma_L = \gamma_D = \gamma_{LD}$ ,  $\varepsilon_L = \varepsilon_D = \varepsilon_{LD}$ , and  $\delta_L = \delta_D = \delta_{LD}$ . In the case of the conformers  $\alpha_L = \alpha_D = \alpha_{LD}$  and  $\varepsilon_L = \varepsilon_D = \varepsilon_{LD}$  the constrained geometry optimizations, performed at the 6-311++G\*\* B3LYP level, utilized characteristic  $\phi$  and  $\psi$  torsional angles (see text for details).

data are presented in Tables V and VI for For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>, respectively. Overall, the agreement between the average theoretical and experimental shifts is rather impressive, especially for



**FIGURE 4.** The 3-21G RHF,  $6-311++G^{**}$  B3LYP,  $6-311++G^{**}$  MP2, and the idealized locations of the different backbone conformers of For-L-Ala-NH<sub>2</sub> on a Ramachandran surface. Only fully optimized geometries are reported.

those data obtained using Levels E and F. It is noteworthy that as the size of the basis set employed increases, the deviation between theoretically and experimentally determined average shifts decreases for all the nuclei but the amide proton. Disregarding the shifts of the amide protons  $({}^{1}H^{NH})$ , the error for all nuclei, both for the Gly and Ala model systems, is down to a few percent. This means, on the absolute (ppm) scale, that the difference between the best calculated and experimentally determined shifts for <sup>13</sup>C and <sup>15</sup>N<sup>NH</sup> is typically in the range of only a few ppm, while for the alpha protons the same difference falls to the tenths of ppm range. At this stage it is not clear to us why the difference between the calculated and observed <sup>1</sup>H<sup>NH</sup> values is so much larger than for any other nucleus. It is perhaps of interest to note that the largest error is associated with the most acidic (i.e., deshielded) protons, H<sup>NH</sup>, and not with the less acidic  $H^{\alpha}$  protons. Though deviation of the calculated  ${}^{1}H^{NH}$  shift from experiment decreases when higher levels of theory are applied, it remains substantial. While the <sup>1</sup>H<sup>NH</sup> shift is off on the absolute (ppm) scale, it is sensitive to conformational changes (vide infra). Until now, theoretical studies were concerned with <sup>13</sup>C and <sup>15</sup>N<sup>NH</sup> shifts only, and with their relation to  $\alpha$ -helical and  $\beta$ -sheet conformers of the peptide models. It is demonstrated here that for all nuclei but one, in both peptide models studied and considering all typical backbone orientations, the average shifts can be determined with considerable accuracy.

#### Chemical Shift—Chemical Shift Correlations for Main-Chain Fold Determination

Typical two-dimensional (2D) single-quantum coherence  $(2D-HSQC)^{11}$  or multiple-quantum coherence  $(2D-HMQC)^{10}$  experiments correlate chemical shifts of selected hetero nuclei (e.g., <sup>13</sup>C and <sup>15</sup>N) with those of their attached protons. It is widely known that not only the side chains modify the observed chemical shifts (e.g., there is a characteristic shift of <sup>15</sup>N<sup>NH</sup> in Gly, Ser, and Thr residues)<sup>22</sup> but for the same residue a significant up- or downfield shift can be detected, depending upon the local molecular environment and/or the backbone conformation. These and similar chemical shift–chemical shift correlations are worthy exploring in detail.

As can be seen from data in Tables III and IV, simple reorientation of the amide planes can have a remarkable effect on the chemical shifts of the nuclei



**FIGURE 5.** <sup>15</sup>N<sup>NH-1</sup>H<sup>NH</sup>-<sup>13</sup>C<sup> $\alpha$ </sup> correlated 3D plot obtained at Level D [**A**] and its three 2D projections (<sup>15</sup>N<sup>NH-1</sup>H<sup>N</sup> [**B**], <sup>13</sup>C<sup> $\alpha$ -1</sup>H<sup>NH</sup> [**C**], and <sup>13</sup>C<sup> $\alpha$ -15</sup>N<sup>NH</sup> [**D**]) for For-L-Ala-NH<sub>2</sub>. All shift values of the nine backbone conformers of For-L-Ala-NH<sub>2</sub> are relative to the energetically most stable conformer,  $\gamma_L$ . NMR computations (6-31+G<sup>\*</sup> GIAO-RHF [square], 6-311++G<sup>\*\*</sup> GIAO-RHF [triangle], and TZ2P GIAO-RHF [circle]) were performed by using two sets of optimized geometries, 6-31++G<sup>\*\*</sup> RHF [empty symbols] and 6-311++G<sup>\*\*</sup> B3LYP [filled symbols].

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**N**<sup>NH</sup> -9

-11

-13

-15

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of our model compounds. For example, for For-L-Ala-NH<sub>2</sub> the  $\gamma_{\rm L} \rightarrow \varepsilon_{\rm D}$  (Level A:  $\phi_{\gamma \rm L} = -84.5^{\circ}$ ,  $\psi_{\gamma \rm L} = 68.7^{\circ} \rightarrow \phi_{\varepsilon \rm D} = 64.7^{\circ}$ ,  $\psi_{\varepsilon \rm D} = -178.6^{\circ}$ ) conformational shift is followed by a 5.5 ppm change in the  $^{13}{\rm C}^{\alpha}$  and a 12.7 ppm change in the  $^{15}{\rm N}^{\rm NH}$  shifts. Furthermore, interconversion of two conformational neighbors, for example that of  $\gamma_{\rm L} \rightarrow \beta_{\rm L}$ , can easily be monitored through changes of selected shifts (e.g., at Level B the appropriate changes are  $\Delta^{15}{\rm N}^{\rm NH} = 9.8$  ppm and  $\Delta^{13}{\rm C}^{\beta} = -4.4$  ppm). These and related effects can clearly be observed on the  $^{15}{\rm N}^{\rm NH}_{-1}{\rm H}^{\rm NH}_{-13}{\rm C}^{\alpha}$  correlated

3D plot [Fig. 5(A)]. Figure 5(A), as well as the the related plots referring to the <sup>15</sup>N<sup>NH</sup>-<sup>1</sup>H<sup>N</sup>, <sup>13</sup>C<sup> $\alpha$ </sup>-<sup>15</sup>N<sup>NH</sup>, and <sup>13</sup>C<sup> $\alpha$ -1</sup>H<sup>HN</sup> planes [Fig. 5(B–D)], reveal distinct locations for the nine backbone conformers of For-L-Ala-NH<sub>2</sub>. Structure-induced changes in the experimentally easily amenable <sup>1</sup>H<sup>NH</sup> shifts are small [ $\Delta\delta$ (<sup>1</sup>H<sup>NH</sup>)  $\approx \pm 0.5$  ppm]. Nevertheless, if the amide proton (<sup>1</sup>H<sup>NH</sup>) shifts are correlated with the amide nitrogen (<sup>15</sup>N<sup>NH</sup>) shifts, a seemingly useful 2D plot emerges [see Fig. 5(B)]: on this <sup>1</sup>H-<sup>15</sup>N HSQC-like plane the characteristic backbone orientations separate clearly. Note, furthermore, that

0,5 H<sup>NH</sup>

1,0

1,5

2.0

all L-type conformers but  $\beta_{\rm L}$  occupy the left part of the 2D plot. Because D-type conformers are rare in globular proteins, the following useful rule emerges for the relative chemical shifts ( $\delta$ -scale in ppm) of <sup>15</sup>N of amides:  $N^{NH}(\gamma_L) > N^{NH}(\delta_L) >$  $N^{\rm NH}(\alpha_{\rm L}) > N^{\rm NH}(\varepsilon_{\rm L}) > N^{\rm NH}(\beta_{\rm L})$ . On the <sup>1</sup>H<sup>NH</sup>-<sup>13</sup>C<sup>\alpha</sup> plane [Fig. 5(C)], the L-type backbone conformations are once again close to each other. In contrast, all D-type conformers are shifted upfield on the <sup>13</sup>C scale. Undoubtedly, each characteristic backbone orientation has its distinct position even on the least commonly used subplane of the 3D correlation plot,  ${}^{13}C^{\alpha}-{}^{15}N^{NH}$  [Fig. 5(D)]. A similarly useful 2D plot of correlated  ${}^{13}C^{\alpha}-{}^{1}H^{\alpha}$  shifts is presented on Figure 6. Overall, in the case of the two model peptides studied, positions of the correlated shifts of the various backbone conformers differ from each other to such an extent that, if this described the real case in peptides and proteins, it would be straightforward to recognize and to assign the backbone conformations of the amino acid residues from their relative chemical shifts. Considering the resolution of modern spectrometers, the chemical shift differences, which might be observed, seem to be sufficient to distinguish and assign the different backbone conformer types of the same residue in a protein.

There are several factors determining the position of the backbone conformers on these correlated plots. At least three of them are related to computa-

**C**<sup>*α*</sup>-**H**<sup>*α*</sup>



**FIGURE 6.** <sup>13</sup>C<sup> $\alpha$ </sup>-<sup>1</sup>H<sup> $\alpha$ </sup> correlated 2D plot. All shifts of the nine characteristic backbone conformers of For-L-Ala-NH<sub>2</sub> are relative to  $\gamma_L$ . NMR computations (6-31+G<sup>\*</sup> GIAO-RHF [square], 6-311++G<sup>\*\*</sup> GIAO-RHF [triangle], and TZ2P GIAO-RHF [circle]) were performed by using two sets of geometries: 6-31++G<sup>\*\*</sup> RHF [empty symbols] and 6-311++G<sup>\*\*</sup> B3LYP [filled symbols].

tions. These include the effects of one-particle basis set deficiency, the extent of electron correlation, and the effect of geometry optimization. Results reported in Tables III and IV as well as data shown on Figures 5 and 6, reveal that the conformers of the two model peptides are clearly distinguishable on the correlated plots, independent of the basis set employed [6-31+G<sup>\*</sup>, 6-311++G<sup>\*\*</sup>, DZ(d) or TZ2P] for the calculation of the chemical shieldings. Although the choice whether the GIAO-RHF calculations were performed at constrained (Levels A and D) or fully optimized (Levels B, C, E, and F) reference geometries of For-L-Ala-NH<sub>2</sub> has an effect on the calculated shieldings (Table IV), separation of the backbone conformers is clearly not affected (Figs. 5–7). In the case of For-Gly-NH<sub>2</sub>, the effect of electron correlation on CSA was determined directly by performing GIAO-RHF and GIAO-MP2 calculations (Table III). The changes in the CSA values of all nuclei (of all conformers) following inclusion of electron correlation (SCF  $\leftrightarrow$  MP2) and extension of basis set  $[DZ(d) \leftrightarrow TZ2P]$  (Levels A and B) have a similar magnitude. More specifically, for the  $\alpha_{LD}$  and  $\varepsilon_{LD}$  conformers of For-Gly-NH<sub>2</sub> the effect of basis set extension slightly overwhelms the effect of electron correlation, while for  $\beta_{\rm LD}$  and  $\delta_{LD}$  the opposite holds. In the future, the effect of electron correlation on the calculated chemical shieldings should be further investigated on chiral peptide models by using much more demanding GIAO-CCSD<sup>47</sup> calculations. Although MP2 calculations have induced significant relative chemical shift changes, it is not expected that electron correlation will alter profoundly the qualitative picture: the basic peptide conformers appear at different regions on chemical shift-chemical shift plots (cf. Figs. 5 and 6). Therefore, the usefulness of the approach of direct determination of conformations of protein building units from multidimensional NMR experiments depends on what effect the side chains, solvation, anisotropic factors, and inter- and intramolecular H-bonding might have on the relative chemical shifts of the selected nuclei. Detailed theoretical investigations of more model compounds and more correlated chemical shielding plots are needed to establish the magnitude of these effects, while it is also hoped that the present theoretical results will encourage experimental work in this direction.

#### **Chemical Shift–Structure Correlations**

One of the primary goals of the research reported in this article has been the investigation of

the use of chemical shifts for predicting the main chain fold of peptides and proteins. So far, we have demonstrated that (a) there is convergence between theoretically determined and experimentally based conformational-average shifts (see above); and (b) chemical shifts cluster as a function of mainchain fold (see previous section). Prompted by these promising results, we performed a systematic linear correlation analysis of all important conformational parameters ( $\phi$ ,  $\psi$ ,  $\omega_0$ ,  $\omega_1$ , and  $\chi_1$ ) and all chemical shifts (<sup>15</sup>N<sup>NH</sup>, <sup>13</sup>C<sup> $\alpha$ </sup>, <sup>13</sup>C<sup> $\beta$ </sup>, <sup>13</sup>C<sup> $\prime$ </sup>, <sup>1</sup>H<sup>NH</sup>, and <sup>1</sup>H<sup> $\alpha$ </sup>). We repeated these linear correlation studies at all 3(6) different levels of theory (see Table I) available for For-Gly-NH<sub>2</sub>(For-L-Ala-NH<sub>2</sub>). This analysis utilizes the linear correlation (Pearson) coefficient, *R*, and the related standard error, *S*<sub>Y,X</sub>, defined as follows:

$$R = \frac{[N(\sum XY) - (\sum X)(\sum Y)]}{\sqrt{[N \sum X^2 - (\sum X)^2][N \sum Y^2 - (\sum Y)^2]}}$$
(1)

and

$$S_{Y.X} = \left( \left[ \frac{1}{N(N-2)} \right] \left[ N \sum Y^{2} - \left( \sum Y \right)^{2} - \frac{\left[ N \sum XY - (\sum X)(\sum Y) \right]^{2}}{N \sum X^{2} - (\sum X)^{2}} \right] \right)^{1/2}$$
(2)

where the vectors **X** and **Y** of length *N* contain the results to be correlated.

A typical set of results ( $R^2$  values), obtained at Level D, is shown in Table VII. For structure–shift correlation only those  $R^2$  values are relevant, which relate structural parameters with shift values, as

given in the box of broken lines in Table VII. (All other  $R^2$  data, reported in Table VII in italics, are irrelevant for this part of the study.) When correlating a set of chemical shifts  $[{}^{13}C^{\alpha}(\alpha_{\rm L}), {}^{13}C^{\alpha}(\beta_{\rm L}),$  etc.] with conformational variables  $[\phi(\alpha_L), \phi(\beta_L), \text{ etc.}]$ , a problem arises due to the fact that torsion angles are variables of periodic nature. In conventional definitions torsions sample regions of  $[-180^\circ, +180^\circ]$  or  $[0^{\circ}, +360^{\circ}]$ . Obviously, the correlation is dependent on which definition is actually applied. To obtain reproducible  $R^2$  values, unless otherwise noted, we selected the period  $[0^\circ, +360^\circ]$  for the definition of the torsion angles. It is further noted that the linear correlation can be optimized by a (moving) shift in the period. The optimum correlation, for each pair, will be described, hereafter, by  $R_{\text{max}}^2$ .

The arithmetical average of the calculated  $R^2$  values is very low; for example, it is 0.14 for the data in Table VII. This indicates that on the average there is no meaningful linear correlation between conformational and shift parameters. However, for certain pairs the  $R^2$  values can be three to four times higher than the average value (e.g.,  $R^2[\phi/C^{\alpha}] =$ 0.488,  $R^2[\phi/H^{\alpha}] = 0.678$ ,  $R^2[\psi/C^{\alpha}] = 0.38$ , and  $R^{2}[\chi_{1}/H^{\alpha}] = 0.371$ ; all printed in bold in Table VII), revealing a significant correlation. The "optimized" correlations,  $R_{max}^2$  (Table VIII), become even more significant. The interesting structure-shift pairs, for example,  $\phi/C^{\alpha}$ , are collected into Table IX for both model systems. The most important conclusions that can be drawn from data in Tables VII-IX are as follows: (a) there is significant correlation be-

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	$\omega_0$	$\phi$	$\psi$	$\omega_1$	χ1	$^{13}C^{\alpha}$	<sup>15</sup> N <sup>NH</sup>	<sup>13</sup> C′	<sup>1</sup> H <sup>NH</sup>	$^{13}C^{\beta}$	<sup>1</sup> Η <sup>α</sup>
ω	1	0.013	0.007	0.636	0.312	0.016	0.001	0.296	0.021	0.000	0.000
$\phi$		1	0.001	0.019	0.254	0.488	0.462	0.278	0.016	0.111	0.678
$\psi$			1	0.002	0.046	0.380	0.086	0.247	0.001	0.059	0.040
$\omega_1$				1	0.139	0.060	0.089	0.022	0.016	0.015	0.052
χ1					1	0.002	0.138	0.250	0.000	0.061	0.371
<sup>13</sup> C <sup>α</sup>						1	0.296	0.000	0.006	0.052	0.452
<sup>15</sup> N <sup>NH</sup>							1	0.003	0.100	0.099	0.423
<sup>13</sup> C′								1	0.007	0.325	0.144
<sup>1</sup> H <sup>NH</sup>									1	0.085	0.019
$^{13}C^{\beta}$										1	0.211
$^{1}$ H $^{\alpha}$											1

<sup>a</sup> All conformational variables are defined in the range [0°, 360°].

<sup>b</sup> Autocorrelation values (e.g.,  $R^2[\phi/\phi]$ ) are all equal to 1, the average of  $R^2$  values is 0.144, with a standard deviation of 0.176. Isotropic shielding and (constrained) geometry optimizations have been performed at the 6-311++G<sup>\*\*</sup> GIAO-RHF and 6-31++G<sup>\*\*</sup> RHF levels, respectively (Level D of Table I).

Maximiz of For-L·	ed R <sup>2</sup> Values (R <sup>2</sup> ma) -Ala-NH <sub>2</sub> . <sup>a</sup>	<sub>x</sub> ) between Selected	d Conformational F	Parameters and Ch	emical Shifts	
	<sup>13</sup> C <sup>α</sup>	<sup>15</sup> N <sup>NH</sup>	<sup>13</sup> C′	<sup>1</sup> H <sup>NH</sup>	<sup>13</sup> C <sup>β</sup>	<sup>1</sup> Η <sup>α</sup>
$\phi$	0.727	0.626	0.340	0.340	0.412	0.678
$\psi$	0.421	0.475	0.247	0.243	0.411	0.264

# TABLE VIII.

<sup>a</sup> See footnotes of Table VII. See text for definition of  $R_{max}^2$ .

#### TABLE IX. \_

#### Correlation of Selected Conformational Parameters ([0°, 360°] Representation) with Chemical Shifts for For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>.<sup>a</sup>

Model	Level	$\phi$ - <sup>1</sup> H <sup><math>\alpha</math>1</sup>	$\phi$ - <sup>1</sup> H <sup><math>\alpha</math>2</sup>	$\phi$ - <sup>13</sup> C $^{\alpha}$	$\psi$ - <sup>13</sup> C $^{\alpha}$	$\psi$ - <sup>15</sup> N <sup>NH</sup>	χ <sub>1</sub> - <sup>13</sup> C′
For-Gly-NH <sub>2</sub>	A B	0.93 (4.08) 0.98 (4.24)	0.49 (2.14) -0.67 (2.00)	0.01 (0.06) -0.22 (0.22)	0.12 (0.55) 0.41 (0.75)	0.21 (0.92) 0.42 (0.78)	
- or-∟-Ala-NH₂	С	0.96 (3.54)	-0.98 (3.67)	-0.59 (1.32)	0.17 (0.11)	-0.02 (0.00)	
For-L-Ala-NH <sub>2</sub>	A B C	-0.83 (4.72) -0.90 (3.57) -0.93 (3.85)		0.68 (3.16) 0.86 (3.30) 0.81 (2.94)	-0.64 (2.80) -0.61 (1.64) -0.71 (2.27)	-0.67 (3.07) -0.51 (1.16) -0.59 (1.57)	0.52 (1.86) 0.53 (1.25) 0.60 (1.60)
	DE	-0.82 (4.72) -0.89 (3.39) 0.70 (3.56)		0.70 (3.39) 0.86 (3.15)	-0.62 (2.65) -0.61 (1.63)	-0.68 (3.21) -0.54 (1.24) 0.55 (1.25)	0.50 (1.74) 0.46 (0.93)
	Г	-0.70 (3.56)		0.03 (3.23)	-0.53 (1.68)	-0.55 (1.25)	0.03 (1.30)

<sup>a</sup> For each conformational parameter–chemical shift correlation first *R*, the linear correlation coefficient [see eq. (1) of the text] is reported followed by  $R^2/\rho$ , where  $\rho$  is the standard deviation of all  $R^2$  values of Table VII. For For-Gly-NH<sub>2</sub> the  $\rho$  values are 0.23, 0.23, and 0.26 for Levels A, B, and C, respectively. For For-L-Ala-NH<sub>2</sub> the p values are {0.14, 0.23, 0.22, 0.14, 0.23, 0.23} for Levels {A, B, C, D, E, F}, respectively.

Nucleus	Conformer	Level A	Level B	Level C	Level D	Level E	Level F
Cα	$\alpha_{D}$	47.10	51.33	47.51	49.40	53.63	53.01
	$\alpha_{L}^{-}$	45.94			47.80		
	$\beta_{L}$	43.36	45.35	43.01	45.31	47.26	46.55
	γD	52.42	55.19	52.38	54.91	57.74	57.28
	γL	41.83	44.84	41.71	43.66	46.72	46.14
	$\delta_{D}$	49.91	51.50	49.09	52.38	54.04	53.44
	$\delta_{L}$	42.13	45.15	42.65	44.07	47.15	46.54
	£р	47.32			49.84		
	$\varepsilon_{L}$	43.19			45.02		
H <sup>α</sup>	$\alpha_{D}$	3.08	3.41	3.07	2.97	3.33	3.28
	$\alpha_{L}$	3.80			3.74		
	$\beta_{L}$	3.77	4.19	3.87	3.74	4.18	4.18
	γd	3.24	3.56	3.24	3.11	3.45	3.42
	γL	3.97	4.29	3.98	3.91	4.23	4.25
	$\delta_{D}$	3.85	4.30	3.94	3.77	4.25	4.29
	$\delta_{L}$	4.51	4.79	4.39	4.53	4.83	4.83
	ε <sub>D</sub>	3.04			2.91		
	εL	3.95			3.91		



**FIGURE 7.** Correlation between *ab initio* determined and experimentally found average (conformation independent) chemical shifts plotted for all relevant nuclei in For-L-Ala-NH<sub>2</sub> at two levels of theory: [A] = GIAO-RHF/6-31+G\*//RHF/6-31++G\*\* (Level A) and [B] = GIAO-RHF/TZ2P//B3LYP/6-311++G\*\* (Level F).

tween  $\phi$  and the <sup>1</sup>H<sup> $\alpha$ </sup> shift, with a typical  $R[\phi/H^{\alpha}]$  value larger than 0.8; (b) the <sup>13</sup>C<sup> $\alpha$ </sup> shift of For-L-Ala-NH<sub>2</sub> correlates with  $\phi$  just as well as with  $\psi$ ; (c) somewhat disturbingly, the  $\phi$ -<sup>13</sup>C<sup> $\alpha$ </sup> and  $\psi$ -<sup>13</sup>C<sup> $\alpha$ </sup> correlations become, at certain levels of theory, almost insignificant for For-Gly-NH<sub>2</sub>, but perhaps this only indicates that achiral For-Gly-NH<sub>2</sub> is not an ideal model for peptides; and (d) none of the correlations seem to become more significant when higher levels of theory are applied for the computation of the underlying quantities.

Although the linear correlation of selected chemical shifts with conformational parameters is not as high as one would hope, pairs with large  $R^2$  values certainly warrant further investigation. All available



**FIGURE 8.** The  $C^{\alpha}/H^{\alpha}$  correlation plot of the nine typical backbone conformers associated with For-L-Ala-NH<sub>2</sub>: (**A**) purely theoretical (Level F), (**B**) corrected with the difference found between calculated and experimental average chemical shifts.

results suggest that relative changes in the chemical shifts of  ${}^{13}C^{\alpha}$  and  ${}^{1}H^{\alpha}$  with the conformational characteristics are important. Therefore, we decided to investigate the dependence of the calculated shifts of these two nuclei as a function of the applied level of theory (Tables X and XI). All chemical shifts calculated at different levels of theory (Levels A-F) correlate well with themselves (e.g.,  $R(^{13}C^{\alpha}) > 0.978$ ,  $R(^{1}H^{\alpha}) > 0.987$ ). This means that extension of the basis or inclusion of electron correlation into the shielding calculation do not increase significantly the reliability of predicting  $[\phi, \psi]$  values from calculated  ${}^{13}C^{\alpha}$  and  ${}^{1}H^{\alpha}$  values. (These findings are based also on results obtained for For-Gly-NH2 but not presented here.) Figure 8(A) reports an ab initio (Level F)  $C^{\alpha}/H^{\alpha}$  correlation plot with respect to all typical backbone conformers of For-L-Ala-NH<sub>2</sub>. Figure 8(B) shows how these ab initio determined shift values would change if a constant, correcting for the

Nucleus	Parameter		А	В	С	D	Е	F					
Cα	R	А	1.000	0.978	0.994	0.999	0.980	0.980					
		В	0.978	1.000	0.992	0.978	1.000	1.000					
		С	0.994	0.992	1.000	0.994	0.993	0.993					
		D	0.999	0.978	0.994	1.000	0.981	0.980					
		E	0.980	1.000	0.993	0.981	1.000	1.000					
		F	0.980	1.000	0.993	0.980	1.000	1.000					
	S <sub>Y.X</sub>	А	_	1.029	0.539	0.156	0.977	0.988					
		В	1.020	_	0.628	1.010	0.098	0.106					
		С	0.523	0.613		0.518	0.569	0.572					
		D	0.168	1.084	0.569		1.024	1.039					
		E	1.034	0.105	0.621	1.018		0.066					
		F	1.058	0.115	0.632	1.045	0.067	—					
Ηα	R	А	1.000	0.992	0.989	0.999	0.991	0.989					
		В	0.992	1.000	0.998	0.991	0.999	1.000					
		С	0.989	0.998	1.000	0.987	0.996	0.998					
		D	0.999	0.991	0.987	1.000	0.992	0.988					
		E	0.991	0.999	0.996	0.992	1.000	0.999					
		F	0.989	1.000	0.998	0.988	0.999	1.000					
	S <sub>Y,X</sub>	А	_	0.074	0.088	0.019	0.076	0.087					
		В	0.074	_	0.032	0.079	0.029	0.018					
		С	0.084	0.031		0.089	0.051	0.037					
		D	0.021	0.087	0.102		0.079	0.097					
		E	0.082	0.032	0.057	0.078		0.025					
		F	0.098	0.020	0.044	0.100	0.026	—					

#### TABLE XI.

Chemical Shift Correlations (*R* and S<sub>Y,X</sub>), as a Function of the Applied Level of Theory (A–F), for the Nuclei  $C^{\alpha}$  and  $H^{\alpha}$  of For-L-Ala-NH<sub>2</sub>.<sup>a</sup>

<sup>a</sup> For primary entries see Table X. See Table I for the description of the theoretical levels.

deviation between experimental and calculated average shift values, were added to them. Figure 8(B) is perhaps the most meaningful way to present the relevant theoretical results to experimentalists.

Because the two predominant secondary structural elements of proteins are the  $\alpha$ -helix and the  $\beta$ -sheet, accurate description of the NMR characteristics of these two conformational building units has been of primary interest for experimentalists and theoreticians alike.<sup>5, 9, 13, 15, 48</sup> The most important results are as follows: (a) there is a characteristic downfield shift of the  $C^{\alpha}$  and C' chemical shifts associated with the right-handed helix conformation; and (b) there is a significant upfield shift when the peptide is in a  $\beta$ -sheet structure. Nevertheless, there is still a lot to be learned as important atoms, such as protons, are typically excluded from these analysis. Consequently, we collected previous theoretical<sup>15c, 19, 22, 49</sup> and experimental chemical shifts in Table XII and compare them with the present chemical shift data. Note that previous shifts might refer to slightly different model compounds with slightly different torsion angles and the calculations have been performed at somewhat different levels of theory. Nevertheless, for peptide models containing alanine (Table XIIA), all calculated chemical shifts for carbon nuclei ( $C^{\alpha}$ ,  $C^{\beta}$ , and C') are in agreement with literature data. The positive value of  $\Delta\delta$  calculated for C<sup> $\alpha$ </sup> agrees with the expected downfield shift of a helical backbone and the upfield shift characteristic of an extended conformation. The smallest difference between calculated ( $\Delta \delta_{\text{calc}}^{\alpha-\beta} = 3.3 \text{ ppm}$ ) and experimentally determined <sup>13</sup>C<sup> $\alpha$ </sup> shifts ( $\Delta \delta_{\exp}^{\alpha-\beta} \approx 3.0$  ppm) is obtained when Level G is compared to a  $3_{10}$ -helix/ $\beta$ -sheet conversion (data not shown). In the case of For-Gly-NH<sub>2</sub>, the best agreement between theory and experiment for <sup>13</sup>C<sup> $\alpha$ </sup> is obtained at the MP2 level ( $\Delta \delta_{exp}^{\alpha-\beta} \approx$ 2.0 ppm and  $\Delta \delta_{\text{calc}}^{\alpha-\beta} = 1.6$  ppm). The sign of the  $\Delta\delta$  quantity for a helix-sheet conversion should be opposite for  $C^{\beta}$  than for  $C^{\alpha}$ , which is indeed the case. Furthermore, the computed difference, typi-

#### TABLE XII.

Comparison of Theoretical and Experimental NMR Relative Chemical Shift Differences ( $\delta^{\alpha-helix} - \delta^{\beta-sheet}$ ) between  $\alpha$ -Helix- and  $\beta$ -Sheet-Like Conformers of the For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub> Model Systems.

For-Gly-NH <sub>2</sub>	N <sup>NH</sup>	Cα	C	,	H <sup>NH</sup>	Hα	Hα	Hα
IGLO/DZ//SCF/3-21G <sup>e</sup> (ref. 19) RHF/6-311++G(2p2d) (ref. 22)		5.2 4.3		-	-4.6			
GIAO-RHF/DZ(d)// B3LYP/6-311++G** (Level A) GIAO-MP2/DZ(d)//B3LYP/6-311++G** (Level B) GIAO-RHF/TZ2P//B3LYP/6-311++G** (Level C)	-3.4 -2.6 5.7	1.0 1.6 4.0	-4. -2. 0.	5 3 2	-1.8 -1.3 -1.4	-0.7 -0.7 -0.5	0.1 0 0.4	-0.3 -0.4 -0.1
Experimental data from Wishart et al. (ref. 5a) Experimental BMRB <sup>f</sup> (Gly[in $\beta$ -sheet]-Gly[in $\alpha$ -helix])		2.8 0.9						-0.2
For-L-Ala-NH <sub>2</sub>			N <sup>NH</sup>	Cα	$C^{eta}$	C′	H <sup>NH</sup>	H <sup>α</sup>
deMon-NMR/IGLOIII <sup>a</sup> (ref. 15c) GIAO-SCF/TZP <sup>a</sup> (ref. 15c) GIAO-SCF/IGLOIII <sup>a</sup> (ref. 15c) RHF/6-311++G(2p2d) (ref. 22) DFT deMon (ref. 43)			-1.7 -2.2 -1.9 	6.8 4.7 5.0 5.0	0 -1.1 -1.2 	0 -0.6 -0.4 	 	  
$ \begin{array}{l} {\hbox{GIAO-RHF/6-31+G^*//SCF/6-31++G^{**b}}$ (Level A) \\ {\hbox{GIAO-RHF/6-311++G^{**//SCF/6-31++G^{**b}}$ (Level D) } \\ {\hbox{RHF/TZ2P//B3LYP/6-311++G^{**b}}$ (Level G) $ {$3_{10}$-helix on RHF/TZ2P//B3LYP/6-311++G^{**c}}$ (Level G) } \end{array} $	rn}	6.3 6.8 6.1 4.5	2.6 2.5 3.3 5.1	-1.4 -1.4 -2.5 -2.3	0.8 1.0 0 0	-1.6 -1.7 -1.9 -2.0	0.0 0.0 0.0 –0.2	
Experimental data from Kricheldorf and Müller (ref. 48) Experimental data from Wishart et al. (ref. 5a) Experimental data from Spera and Bax (ref. 9d) Experimental BMRB <sup>d</sup> (Ala[in $\beta$ -sheet] – Ala[in $\alpha$ -helix])	1		0	3.5 3.5 4.6 3.0	-4.8  -2.5	4.6 2.1 —	  	  0.8

<sup>a</sup> Computed values are for *N*-acetyl-*N'*-methylalaninamide; see Table I of ref. 15c for results for the  $\alpha$ -helix ( $\phi = -60^{\circ}, \psi = -60^{\circ}$ ) and the extended conformer ( $\phi = -120^{\circ}, \psi = 120^{\circ}$ ).

<sup>b</sup> Values calculated for *N*-formyl-alaninamide are for 3<sub>10</sub>-helices and/or for residue (*i* + 1) in a type I  $\beta$ -turn ( $\phi = -68.6^{\circ}, \psi = -17.5^{\circ}$ ) and extended ( $\phi = -167.6^{\circ}, \psi = 169.9^{\circ}$ ) conformers.

<sup>c</sup> Values calculated for *N*-formyl-alaninamide are for the  $\alpha$ -helix ( $\phi = -54^{\circ}, \psi = -45^{\circ}$ ) and for the extended ( $\phi = -167.6^{\circ}, \psi = 169.9^{\circ}$ ) conformer.

<sup>d</sup> Some 50 alanine residues, all having  $C^{\alpha}$  and  $H^{\alpha}$  assignments, have been included.

<sup>e</sup> Values calculated for *N*-acetyl-*N'*-methylglycinamide are for the  $\alpha$ -helix ( $\phi = -48^{\circ}, \psi = -57^{\circ}$ ) and for the extended ( $\phi = -139^{\circ}, \psi = 135^{\circ}$ ) conformers taken from Table I of ref. 19.

<sup>f</sup> Some 40 glycine residues, all having C<sup> $\alpha$ </sup> and H<sup> $\alpha$ </sup> assignments, have been included.

cally around  $-2.5 \text{ ppm} \leq \Delta \delta_{\text{calc}}^{\alpha-\beta} \leq -1.1 \text{ ppm}$ , agrees well with the experimental value of  $\Delta \delta_{\text{exp}}^{\alpha-\beta} \approx$  $-2.5 \text{ ppm.}^{9\text{c}}$  Carbonyl carbons are less frequently used for monitoring conformational shifts. Once again, our calculations confirm that <sup>13</sup>C  $\sigma$  values contain conformational information<sup>50, 51</sup> and chemical shifts have information about  $\phi, \psi$  values. The result obtained at Level D (Table XIIA) is the closest to the experimental value of Wishart and Sykes,<sup>5a</sup> deviating by only about 1 ppm ( $\Delta \delta_{\text{calc}}^{\alpha-\beta} = 1.0 \text{ ppm}$ ,  $\Delta \delta_{\text{expt}}^{\alpha-\beta} = 2.1 \text{ ppm}$ ). Although <sup>15</sup>N<sup>NH</sup> shifts are also affected by dihedral parameters,<sup>51</sup> they are expected

to be more complicated to use for a conformational analysis. For example, the experimental value for the amide nitrogen,  $\Delta\delta^{\alpha-\beta}_{expt}\approx 0$  ppm, has an uncertainty of  $\pm 4$  ppm.  $^{5a}$  As reported in Table XII, values computed for the  $^{15}N^{NH}$  shift in different conformational states can be rather different.

#### Conclusions

Establishing correlation between peptide mainchain folds and chemical shifts provides a continuous challenge for experimentalists and theoreticians alike. For theoreticians, uncertainties arise, for example, from the problem of ideal size and type of a peptide model to be used for the computations, the required minimum level of ab initio theory, and the incorporation of important structural factors (e.g., nonplanarity of the amide groups). There are just as severe experimental difficulties in establishing such correlations. Consequently, unambiguous correlations have been put forward only for the  $\alpha$ helical and  $\beta$ -sheet regions of the Ramachandran surface. One of the principal aims of this computational study has been the confirmation of existing correlations and derivation of new ones. Some more important findings of this study, relevant or related to this issue, are as follows:

The energy order of the L-type conformers of the two peptide models, For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>, is identical, regardless whether electron correlation is incorporated into the computation or not, and it is as follows:  $E(\gamma_L) < E(\beta_L) < E(\delta_L) < E(\varepsilon_L) < E(\alpha_L)$ . Typically, the D-type backbone conformers have higher relative energy or lower stability. The  $\alpha_L$ ,  $\varepsilon_L$ , and  $\varepsilon_D$  conformers are found not to correspond to local minima on the potential energy hypersurfaces of For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>.

Agreement between theoretical and experimental (BMRB) chemical shifts averaged over all conformers is impressive. As the size of the basis set employed for the shielding calculation increases, the deviation between the shifts decreases to a few percent for all nuclei but the amide proton.

If <sup>1</sup>H<sup>NH</sup> shifts are correlated with <sup>15</sup>N<sup>NH</sup> shifts (the <sup>15</sup>N<sup>H</sup>-<sup>1</sup>H<sup>N</sup> planes could result from HSQC (<sup>15</sup>N<sup>NH</sup>, <sup>1</sup>H<sup>NH</sup>)-type experiments) or if <sup>1</sup>H<sup> $\alpha$ </sup> shifts are plotted against <sup>13</sup>C<sup> $\alpha$ </sup> shifts, useful 2D plots emerge, independently of the basis set employed and whether electron correlation has been included in the theoretical treatment. In the case of the two model peptides studied, positions of the various backbone conformers on the chemical shiftchemical shift plots differ from each other to such an extent that, if this described the real case in peptides, it would be straightforward to recognize and to assign the main-chain fold from the relative chemical shifts.

The empirical structure–chemical shift correlations observed by Wishart,<sup>5</sup> Oldfield,<sup>13</sup> Bax,<sup>4d</sup> and others have been successfully applied in structural determinations of helical and extended-like secondary structures. The data presented in this paper give further justification to these empirical findings and reveals the intrinsic correlation of all nine characteristic backbone conformations of peptides and their chemical shift information.

In summary, ab initio isotropic NMR shielding results presented in this article for the model systems For-Gly-NH2 and For-L-Ala-NH2 facilitate and encourage the application of correlated relative chemical shift information from {<sup>1</sup>H-<sup>15</sup>N}HSQC, {<sup>1</sup>H-<sup>13</sup>C}HSQC, HNCA, HNCB, and other multiplepulse NMR experiments to extract structural information directly from these measurements, thus opening an alternative route to NOE's to derive structures of proteins from their NMR spectra. Detailed theoretical investigations of more model compounds (e.g., For-Ser-NH<sub>2</sub>, For-Val-NH<sub>2</sub>, and For-Phe-NH<sub>2</sub>) are underway in our laboratories. The preliminary results support all the conclusions presented in this article for the simpler model systems For-Gly-NH<sub>2</sub> and For-L-Ala-NH<sub>2</sub>.

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